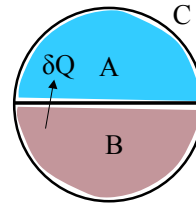


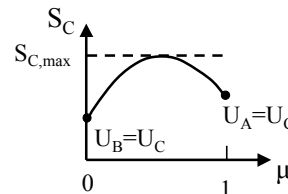
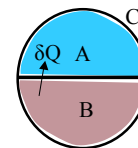
## Entropy Approach to 2<sup>nd</sup> Law for CM

- Start with existence of entropy as a **thermodynamic property** (measure of microscopic disorder of system, or better, energy states of system)
  - $S=S(U,V)$  function of two indep. variables
- Imagine two systems, **A** and **B**, separated by rigid wall ( $V_A, V_B$  constant), only exchange is heat transfer,  $\delta Q$
- Combined system **C** is **isolated**
  - no interaction with surroundings
  - total energy  $U_C$  fixed
  - 2<sup>nd</sup> Law,  $S_{C,final} - S_{C,initial} = P_s \geq 0$



## Equilibrium at Maximum Entropy

- Assume A and B not in equilibrium with each other
- Since entropy is **extensive** property
  - $S_C = S_A(U_A, V_A) + S_B(U_B, V_B)$
  - valid if A, B individually in equilibrium
  - also,  $U_C = U_A + U_B = \mu U_C + (1-\mu)U_C$
- As interaction occurs,  $S_C$  will increase until equilibrium achieved ( $S_{C,max}$ )
- Since  $V_A, V_B$  and  $U_C$  fixed, only free variable is  $\mu$ 
  - maximize  $S_C, dS_C/d\mu = 0$



## Thermodynamic Def'n. of Temperature

- Taking derivative (using chain-rule)

$$\begin{aligned} \frac{dS_C}{d\mu} &= \left( \frac{\partial S_A}{\partial U_A} \right)_{V_A} \frac{dU_A}{d\mu} + \left( \frac{\partial S_B}{\partial U_B} \right)_{V_B} \frac{dU_B}{d\mu} \\ &= \left( \frac{\partial S_A}{\partial U_A} \right)_{V_A} U_C + \left( \frac{\partial S_B}{\partial U_B} \right)_{V_B} (-U_C) \end{aligned}$$

- Equals zero, when  $\left( \frac{\partial S_A}{\partial U_A} \right)_{V_A} = \left( \frac{\partial S_B}{\partial U_B} \right)_{V_B}$
- So, two systems are in thermal equilibrium when they have same  $(\partial S/\partial U)_V$
- But, same T also implies thermal equilibrium
- Therefore, let  $T \equiv \frac{1}{(\partial S/\partial U)_V}$  (show reason for 1/ next)

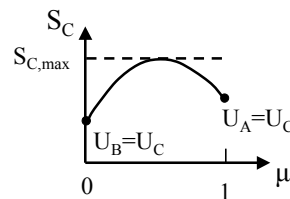
## Thermodynamic Def'n. of T (con't)

- Need to show thermodynamic def'n. of T consistent with 0<sup>th</sup> Law of Thermo. (Q from high to low T)

$$\frac{dS_C}{d\mu} = \left( \frac{\partial S_A}{\partial U_A} \right)_{V_A} U_C + \left( \frac{\partial S_B}{\partial U_B} \right)_{V_B} (-U_C)$$

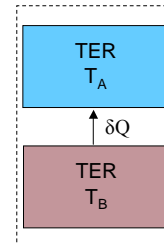
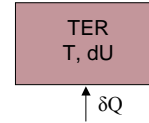
$$dS_C = \left( \frac{1}{T_A} - \frac{1}{T_B} \right) U_C d\mu$$

- If  $T_B > T_A$ ,  $\left( \frac{1}{T_A} - \frac{1}{T_B} \right) > 0$
- 2<sup>nd</sup> Law  $\Rightarrow dS_C > 0$ , so  $d\mu > 0$  energy flowing from B to A
- For  $dS_C/d\mu$  to truly be maximum, then  $d^2S_C/d\mu^2 < 0$   
– result,  $\left( \frac{\partial^2 S}{\partial U^2} \right)_V < 0 \Rightarrow \left( \frac{\partial T}{\partial U} \right)_V > 0$  U increases with T



## Heat Transfer as Entropy Transfer

- Idealize, **Thermal Energy Reservoir (TER)**
  - CM with fixed volume, only exchanges energy as  $Q$ , uniform and (nearly) constant  $T$
  - Since fixed volume,  $dS/dU=1/T$
  - From 1<sup>st</sup> Law,  $dU=\delta Q$ , so  $dS_{TER} = \delta Q/T$
- Now consider two interacting TERs, isolated from surroundings
  - 2<sup>nd</sup> Law,  $\delta P_s = d(S_A + S_B)$   
 $= \delta Q(1/T_A - 1/T_B) > 0$
  - entropy production associated with  $Q$  across finite temperature difference



## Second Law for Control Mass

- Control mass interacting with TER and **Mechanical Energy Reservoir (MER)**
  - CM with no microscopic disorder (no entropy), can only exchange energy as reversible work
- Together, form isolated system
- 2<sup>nd</sup> Law
  - $\delta P_s = d(S_{TER} + S_{MER} + S)$   
 $= -\delta Q/T + dS$
  - in terms of inequality

$$dS = \delta Q/T + \delta P_s$$

$$dS \geq \delta Q/T$$

