

## Radiation

- So far we have used statistical mechanics/statistical thermodynamics to examine equilibrium properties of matter
- Thermodynamics, being the study of matter and energy, is also applicable to pure energy, e.g., electromagnetic radiation
- As with our previous study of matter, the primary question is
  - What is the distribution of the “particles” among the possible energy levels at equilibrium?
- For radiation, the “particles” are **photons**

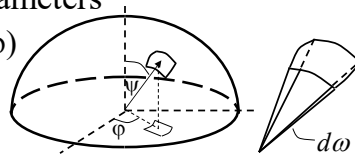


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## Radiation Field Descriptors

- Photons characterized by 2 parameters
  - 1) direction of propagation ( $\psi, \phi$ )
    - differential 2-d angular spread given by  $d\omega \equiv$  **differential solid angle**
  - 2) frequency  $\nu$  (or wavelength  $\lambda = c/\nu$ ,  $c$ =speed of light)
    - photon energy,  $\varepsilon = h\nu$
    - photon momentum,  $p = h/\lambda = h\nu/c$  *NOTE:  $p \neq$  pressure*
- Compare to mass particles, characterized by
  - 1) direction of motion
  - 2) speed ( $\Rightarrow$  energy and momentum)



$$d\omega = \sin\psi \, d\psi \, d\phi$$

(units=steradians or sr)

$$\int d\omega = 4\pi$$

*so photons are different – all have same speed (in vacuum), but different energies and momenta*

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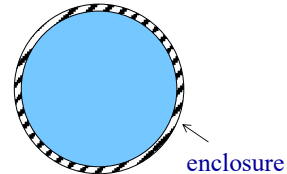
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## Photon Distribution Function

- In general, a radiation field can be described by
  - **intensity**  $I$ , e.g.,  $W/m^2/(sr)$ , or *sr unit often omitted*
  - **photon density**  $n^R$ , e.g., number of photons/ $m^3/(sr)$
- The **photon frequency distribution function** can be written as  $f^R = f^R(\nu, \omega)$  such that
  - $n^R f^R d\omega d\nu$  = number of photons per unit volume in frequency (energy) interval  $\nu \rightarrow \nu + d\nu$  and given center direction within angular “cone”  $d\omega$
- Spectral density** is defined by  $n_\nu^R \equiv n^R f^R$ 
  - number of photons/ $m^3/Hz/(sr)$
- Spectral intensity** is  $I_\nu = c h \nu n_\nu^R$ 
  - $W/m^2/Hz/(sr)$

## Radiative Equilibrium

- What do  $I_\nu$  or  $n_\nu^R$  look like if the radiation energy mode is in “internal” **equilibrium**
  - i.e., equilibrium distribution of photons among the possible energy levels
- To find this, consider a radiation field interacting with matter
  - **radiation field** inside (vacuum) enclosure that is opaque to radiation, i.e., nontransmitting and absorbs all incident radiation
  - since radiation can not escape, will eventually reach equilibrium with enclosure
  - if in equilibrium with radiation field, then body must also emit equilibrium radiation, known as a **blackbody**
- Can think of this as an equilibrium photon gas



## Approach for Equilibrium Distribution

- Use same statistical mechanics approach developed for ideal gas (“particles”)
  - but photon gas follows *Bose-Einstein statistics* (can have unlimited number of bosons in a single quantum state)
- Previously presented that most probable macrostate for *B-E* statistics has distribution  $N_i^* = \frac{g_i}{e^{\alpha+\beta\epsilon_i} - 1}$ 
  - where  $\alpha$  was Lagrange multiplier to satisfy constraint  $N = \sum N_i$ , and  $\beta$  was to satisfy  $E = \sum N_i \epsilon_i$
- For photons, there is no constraint on  $N \Rightarrow \alpha=0$ 
  - nothing prevents 1 high energy photon from becoming 2 low energy photons (no equivalent to mass conservation)
- Using Boltzmann relation ( $S=k\ln\Omega$ ), and TD expression  $(\partial S/\partial E)_{V,N} = 1/T$  again gives  $\beta = 1/kT$

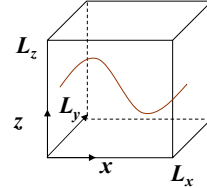
## Equilibrium Frequency Distribution

- So equilibrium (most probable) photon distribution is
 
$$N_i^* = \frac{g_i}{e^{\epsilon_i/kT} - 1}$$
- Assuming any photon frequency (energy) is possible, i.e., the distribution is continuous, and with  $\epsilon_i = h\nu$ 

$$N(\nu) d\nu = \frac{g(\nu)}{e^{h\nu/kT} - 1} \quad \text{from this point, will drop * notation; all results for equilibrium distributions}$$
- This leads to question, **what is degeneracy  $g(\nu)$** ?
- For our continuous distribution, this is equivalent to asking how many quantum states are in the frequency range  $\nu$  to  $\nu + d\nu$ 
  - this is similar to asking about extended degeneracy for translating particles in a box

## Radiation Enclosure

- Examine radiation enclosure using similar approach as traveling (translating) particle in a box; consider eigenvalues for **standing EM waves** in enclosure
- Translating particle in a box quantum numbers were  $n_i = K_i L_i / \pi$ 
  - periodic boundary conditions gave  $K_i = (2\pi/h)(mv_{i,n_i})$
- Recognize that  $mv$  was particle momentum
  - rewriting for photons, with momentum  $= h\nu/c$ , and using isotropy (same result in any direction)
    - $\Rightarrow$  photon quantum numbers in enclosure are  $n = (2L/c)\nu$
    - or  $n^3 = (8V/c^3)\nu^3$



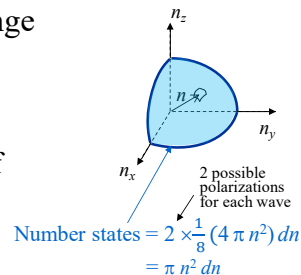
## Degeneracy

$$n^3 = (8V/c^3)\nu^3$$

- Recall the photon degeneracy is number of quantum states in the frequency range  $\nu$  to  $\nu + d\nu$
- This is related to number of quantum states in range  $n$  to  $n + dn$ 
  - which is number of state in portion of spherical shell in “ $n$ -space” of radius  $n$  that is restricted to  $n_x, n_y, n_z > 0$
  - equals  $\pi n^2 dn$
- From our relation between  $n$  and  $\nu$

$$d(n^3) = (8V/c^3) d(\nu^3)$$

$$\Rightarrow \pi n^2 dn = (8\pi V/c^3) \nu^2 d\nu = g(\nu) d\nu$$



## Equilibrium (Blackbody) Radiation

- So inserting this into our earlier equilibrium photon distribution function

$$N(\nu) d\nu = \frac{g(\nu)}{e^{h\nu/kT} - 1} = \frac{8\pi V}{c^3} \frac{\nu^2 d\nu}{e^{h\nu/kT} - 1}$$

- In terms of the photon spectral density

$$n_\nu^R = \frac{N(\nu)}{4\pi V} = \frac{2\nu^2/c^3}{e^{h\nu/kT} - 1}$$

*4π steradians (isotropic)*

- In terms of spectral intensity

$$I_{\nu,b} = h\nu c n_\nu^R = \frac{2h\nu^3/c^2}{e^{h\nu/kT} - 1}$$

*Planck added this after postulating quantized radiated energy, so able to predict long λ radiation behavior*

- Planck Function
- Planck's Law
- Planck Blackbody Distribution, so often denoted  $I_{\nu,b}$

- in terms of wavelength

$$I_\nu = dI/d\nu,$$

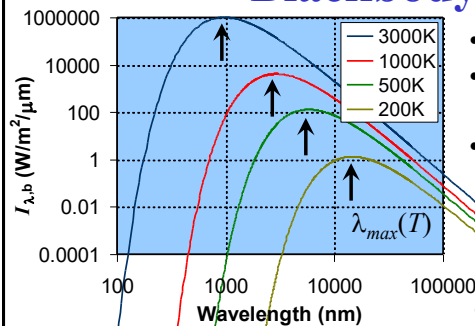
$$\lambda = c/\nu, d\lambda/d\nu = -c/\nu^2$$

$$I_{\lambda,b} = \frac{2hc^2/\lambda^5}{e^{hc/\lambda kT} - 1}$$

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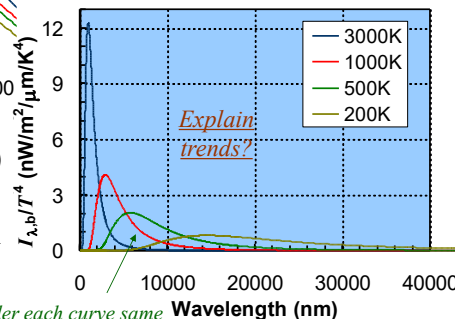
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## Blackbody Intensity



- $I_{\lambda,b} \uparrow$  with  $T$  for all  $\lambda$
- $\lambda_{max}$  shifts blue as  $T \uparrow$   
–  $\lambda_{max} T = 2897.8 \mu\text{m K}$
- $I_{\lambda_{max},b} / T^5 = \text{constant}$   
 $\sim 0.004 \text{ nW/m}^2/\mu\text{m/K}^5$

- Spectrally integrated  $I_b = \text{fn}(T)$   
*Stefan-Boltzmann constant*  
–  $I_b/T^4 \sim 18 \text{ nW/m}^2/\text{K}^4 \equiv \sigma/\pi$   
– so total radiation intensity (energy) scales as  $T^4$



*Area under each curve same*

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## Photon Distribution Function

- In terms of the distribution function

$$f^R(\nu) = \frac{I_\nu}{I_{tot}} = \frac{n_\nu^R}{n^R} = \frac{\pi}{\sigma T^4} \frac{2h\nu^3/c^2}{e^{h\nu/kT} - 1}$$

- So like matter, we can describe a radiation field, **that is in equilibrium**, with a single parameter, its temperature
  - the temperature is a measure of how the “particles” (photons) are distributed among the energy levels
- Unlike many of the energy modes of matter, radiation is often far from being in equilibrium
  - especially the radiation fields produced by gases at typical pressures

## Equilibrium Radiation and Energy

- Examine total equilibrium radiation intensity as function of temperature,  $I_b(T)$
- What is the total energy of an atomic gas (neglecting electronic energy)?
  - $E = 3/2 NkT$   
 $= 3/2 pV$
- How long to deplete all this energy due to blackbody radiation?
  - $t = 3/2 p(V/A)/I_b$
  - for 1mm diameter spherical region at 1bar

